

Time-dependent Hamiltonian $H = \frac{1}{2} \hbar \omega(t) \hat{n}(t) \cdot \vec{\sigma}$

Short-time approximation $T(t+dt, t) = e^{-i(\omega(t)dt/2)} \hat{n}(t) \cdot \vec{\sigma}$

Exactly solvable case $\omega(t)$ constant $\hat{n}(t)$ rotating (NMR)

$$\omega \hat{n}(t) = \omega_z \hat{z} + \omega_p (\cos \omega_r t \hat{x} + \sin \omega_r t \hat{y})$$

" ω_z " ω_p "pulse" "rotating" radio wave
 " " amplitude frequency ω_r

The diagram shows a central purple sphere labeled "sample" positioned within a blue vertical cylinder labeled "B_z". A horizontal yellow arrow labeled "transmitter pulse" points upwards from the left. A red arrow labeled "detector pickup coil" points downwards from the right.

Trick for solving - rotating coordinate system

Given $|\psi\rangle$ in basis $\{|b_j\rangle\}$: $|\psi\rangle = \left(\sum_j |b_j\rangle \langle b_j| \right) |\psi\rangle = \sum_j (\langle b_j | \psi \rangle) |b_j\rangle$

express in basis $\{|\bar{b}_k\rangle\}$: $\langle \bar{b}_k | \psi \rangle = \langle \bar{b}_k | \left(\sum_j |b_j\rangle \langle b_j| \right) |\psi\rangle = \sum_j \langle b_j | \psi \rangle$

Transformation of Coordinate System

Unitary matrix $\langle \bar{b}_k | b_j \rangle$

$$|\psi\rangle \rightarrow |\bar{\psi}\rangle = S(t)|\psi\rangle \text{ inverse } |\psi\rangle = S^+|\bar{\psi}\rangle$$

original basis new basis unitary

$$\begin{pmatrix} |\psi_1\rangle \\ |\psi_2\rangle \\ \vdots \\ |\psi_N\rangle \end{pmatrix} \quad \begin{pmatrix} |\bar{\psi}_1\rangle \\ |\bar{\psi}_2\rangle \\ \vdots \\ |\bar{\psi}_N\rangle \end{pmatrix}$$

$$\text{Schrödinger} \Rightarrow i\hbar \frac{d}{dt} |\bar{\psi}\rangle = i\hbar \frac{dS}{dt} |\psi\rangle + S \left(i\hbar \frac{d}{dt} |\psi\rangle \right)$$

$$= (i\hbar \dot{S} S^+ + S H S^+) |\bar{\psi}\rangle$$

\bar{H}_1 new term when $S \neq 0$ \bar{H}_0 original H in new coords

$$\text{Set } H = \frac{1}{2} \omega_z \hat{\sigma}_z + \frac{1}{2} \omega_p (\cos(\omega_r t) \hat{\sigma}_x + i \sin(\omega_r t) \hat{\sigma}_y)$$

Note: $\hbar = 1$

$$\text{Let } S(t) = \begin{pmatrix} e^{i\omega_r t/2} & 0 \\ 0 & e^{-i\omega_r t/2} \end{pmatrix}$$

Frame rotates around \hat{z} at frequency ω_r ... matches RF

$$\bar{H}_1 = i \dot{S} S^+ = i \begin{pmatrix} \frac{i\omega_r}{2} e^{i\omega_r t/2} & 0 \\ 0 & -\frac{i\omega_r}{2} e^{-i\omega_r t/2} \end{pmatrix} \begin{pmatrix} e^{-i\omega_r t/2} & 0 \\ 0 & e^{+i\omega_r t/2} \end{pmatrix}$$

$$= -\frac{1}{2} \omega_r \hat{\sigma}_z$$

Effective magnetic field $\parallel -\hat{z}$

$$\bar{H}_0 = S H S^+ = \frac{1}{2} \omega_z \hat{\sigma}_z + \frac{1}{2} \omega_p \begin{pmatrix} e^{i\omega_r t/2} & 0 \\ 0 & e^{-i\omega_r t/2} \end{pmatrix} \begin{pmatrix} 0 & c-is \\ c+is & 0 \end{pmatrix} \begin{pmatrix} e^{-i\omega_r t/2} & 0 \\ 0 & e^{+i\omega_r t/2} \end{pmatrix}$$

$$= \frac{1}{2} \omega_z \hat{\sigma}_z + \frac{1}{2} \omega_p \hat{\sigma}_x$$

In frame rotating at ω_r ,
 $\vec{B}_{xy} \rightarrow B_x$ constant in time!

≈ 0 if $\omega_r = \omega_z$

$$\bar{H} = \bar{H}_0 + \bar{H}_1 = \frac{1}{2} (\omega_z - \omega_r) \hat{\sigma}_z + \frac{1}{2} \omega_p \hat{\sigma}_x$$

On resonance $\bar{H} = \frac{1}{2} \omega_p \hat{\sigma}_x$

$$T = e^{-i(\omega_p t/2) \hat{\sigma}_x}$$

$$= \cos\left(\frac{\omega_p t}{2}\right) I - i \sin\left(\frac{\omega_p t}{2}\right) \hat{\sigma}_x$$

| t | T | $ \psi\rangle$ |
|-----|-----|----------------------|
| 0 | I | $ \tilde{z}+\rangle$ |

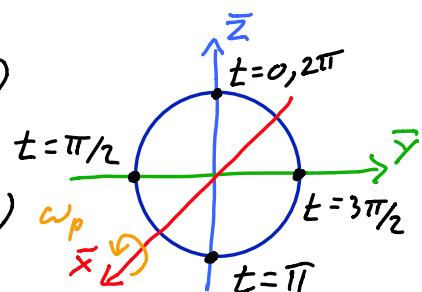
$$\pi/2 \quad \frac{1}{\sqrt{2}} (I - i \hat{\sigma}_x) \quad |\tilde{y}-\rangle = \frac{1}{\sqrt{2}} (|\tilde{z}+\rangle - i |\tilde{z}-\rangle)$$

$$\pi \quad -i \hat{\sigma}_x \quad -i |\tilde{z}-\rangle$$

$$3\pi/2 \quad -\frac{1}{\sqrt{2}} (I + i \hat{\sigma}_x) \quad |\tilde{y}+\rangle = \frac{-1}{\sqrt{2}} (|\tilde{z}+\rangle + i |\tilde{z}-\rangle)$$

$$2\pi \quad -I \quad -|\tilde{z}+\rangle$$

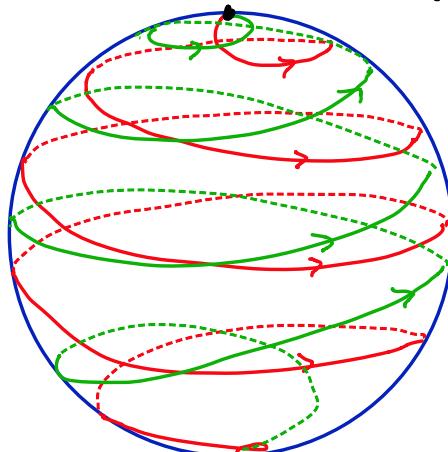
units $1/\omega_p$



Evolution of $\vec{S}(t)$

one full period $\frac{2\pi}{\omega_p}$

π pulse $t = \frac{\pi}{\omega_p}$



$$|\psi_0\rangle = |z+\rangle, |\psi_{\pi}\rangle = -|z+\rangle$$

x-y rotation \curvearrowleft
at frequency ω_r

Z rotation at
frequency ω_p

$$|\psi_{\pi}\rangle = -i|z-\rangle$$

Off Resonance Define $\delta = \omega_z - \omega_r \Rightarrow \bar{H} = \frac{1}{2} \begin{pmatrix} \delta & \omega_p \\ \omega_p & -\delta \end{pmatrix}$

$$\text{Eigenvalues } \lambda = \pm \frac{1}{2} \Omega, \Omega = \sqrt{\omega_p^2 + \delta^2}$$

$$\text{Eigenvectors } |\bar{\varphi}_{\pm}\rangle = A \begin{pmatrix} \delta \pm \Omega \\ -\omega_p \end{pmatrix} \quad A_{\pm} = \frac{1}{\sqrt{\omega_p^2 + (\delta \pm \Omega)^2}}$$

$$T = e^{-it\bar{H}} = e^{-i\Omega t/2} |\bar{\varphi}_+\rangle \langle \bar{\varphi}_+| + e^{i\Omega t/2} |\bar{\varphi}_-\rangle \langle \bar{\varphi}_-|$$

Born Rule $P_r([\bar{z}-] | \psi_0 = |z+\rangle) = |\langle z- | T(t) | z+\rangle|^2$

$$= \frac{\omega_p^2}{\omega_p^2 + (\omega_r - \omega_z)^2} \sin^2\left(\frac{\Omega t}{2}\right)$$

