

Time-dependent Hamiltonian $H = \frac{1}{2} \hbar \omega(t) \hat{n}(t) \cdot \vec{\sigma}$

Short-time approximation $T(t+\Delta t, t) = e^{-i(\omega(t)\Delta t/2) \hat{n}(t) \cdot \vec{\sigma}}$

Exactly solvable case $\omega(t)$ constant $\hat{n}(t)$ rotating (NMR)

$$\omega \hat{n}(t) = \underbrace{\omega_2}_{B_2} \hat{z} + \underbrace{\omega_p}_{\text{"pulse" amplitude}} \left(\underbrace{\cos \omega_r t}_{\text{"rotating" frequency } \omega_r} \hat{x} + \sin \omega_r t \hat{y} \right)$$

Trick for solving - rotating coordinate system

Given $|\psi\rangle$ in basis $\{|b_j\rangle\}$: $|\psi\rangle = \left(\sum_j |b_j\rangle \langle b_j| \right) |\psi\rangle = \sum_j \langle b_j | \psi \rangle |b_j\rangle$

express in basis $\{|\bar{b}_k\rangle\}$: $\langle \bar{b}_k | \psi \rangle = \langle \bar{b}_k | \left(\sum_j |b_j\rangle \langle b_j| \right) |\psi\rangle = U_{kj} \langle b_j | \psi \rangle$

Transformation of Coordinate system

Unitary matrix $\langle \bar{b}_k | b_j \rangle$

$$|\psi\rangle \rightarrow |\bar{\psi}\rangle = S(t) |\psi\rangle \quad \text{inverse } |\psi\rangle = S^\dagger |\bar{\psi}\rangle$$

original basis $\begin{pmatrix} \psi_1 \\ \psi_2 \\ \vdots \\ \psi_N \end{pmatrix}$ new basis $\begin{pmatrix} \bar{\psi}_1 \\ \bar{\psi}_2 \\ \vdots \\ \bar{\psi}_N \end{pmatrix}$ unitary

$$\text{Schrödinger} \Rightarrow i\hbar \frac{d}{dt} |\bar{\psi}\rangle = i\hbar \frac{dS}{dt} |\psi\rangle + S \left(i\hbar \frac{d}{dt} |\psi\rangle \right)$$

$$= (i\hbar \dot{S} S^\dagger + S H S^\dagger) |\bar{\psi}\rangle$$

\bar{H}_1 new term when $\dot{S} \neq 0$ \bar{H}_0 Original H in new Coords

Set $H = \frac{1}{2} \omega_z \sigma_z + \frac{1}{2} \omega_p (\cos(\omega_r t) \sigma_x + i \sin(\omega_r t) \sigma_y)$

Note: $\hbar = 1$

Let $S(t) = \begin{pmatrix} e^{i\omega_r t/2} & 0 \\ 0 & e^{-i\omega_r t/2} \end{pmatrix}$ Frame rotates around \hat{z} at frequency $\omega_r \dots$ matches RF

$$\bar{H}_1 = i \dot{S} S^\dagger = i \begin{pmatrix} \frac{i\omega_r}{2} e^{i\omega_r t/2} & 0 \\ 0 & -\frac{i\omega_r}{2} e^{-i\omega_r t/2} \end{pmatrix} \begin{pmatrix} e^{-i\omega_r t/2} & 0 \\ 0 & e^{i\omega_r t/2} \end{pmatrix}$$

$= -\frac{1}{2} \omega_r \sigma_z$ effective magnetic field $\parallel -\hat{z}$

S commutes with σ_z

$$\bar{H}_0 = S H S^\dagger = \frac{1}{2} \omega_z \sigma_z + \frac{1}{2} \omega_p \begin{pmatrix} e^{i\omega_r t/2} & 0 \\ 0 & e^{-i\omega_r t/2} \end{pmatrix} \begin{pmatrix} 0 & c-is \\ c+is & 0 \end{pmatrix} \begin{pmatrix} e^{-i\omega_r t/2} & 0 \\ 0 & e^{i\omega_r t/2} \end{pmatrix}$$

$= \frac{1}{2} \omega_z \sigma_z + \frac{1}{2} \omega_p \sigma_x$ In frame rotating at ω_r , $\vec{B}_{xy} \rightarrow B_x$ constant in time!

$\neq 0$ if $\omega_r = \omega_z$

$$\bar{H} = \bar{H}_0 + \bar{H}_1 = \frac{1}{2} (\omega_z - \omega_r) \sigma_z + \frac{1}{2} \omega_p \sigma_x$$

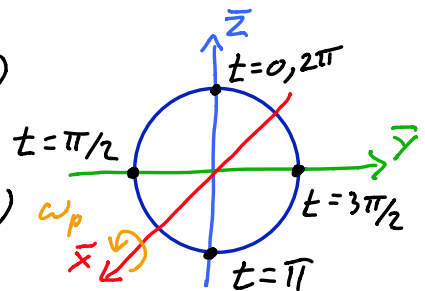
On resonance $\bar{H} = \frac{1}{2} \omega_p \sigma_x$

$$T = e^{-i(\omega_p t/2) \sigma_x}$$

$$= \cos\left(\frac{\omega_p t}{2}\right) I - i \sin\left(\frac{\omega_p t}{2}\right) \sigma_x$$

| t | T | $ \psi\rangle$ |
|----------|---------------------------------------|--|
| 0 | I | $ \bar{z}+\rangle$ |
| $\pi/2$ | $\frac{1}{\sqrt{2}} (I - i\sigma_x)$ | $ \bar{y}-\rangle = \frac{1}{\sqrt{2}} (\bar{z}+\rangle - i \bar{z}-\rangle)$ |
| π | $-i\sigma_x$ | $-i \bar{z}-\rangle$ |
| $3\pi/2$ | $-\frac{1}{\sqrt{2}} (I + i\sigma_x)$ | $ \bar{y}+\rangle = \frac{1}{\sqrt{2}} (\bar{z}+\rangle + i \bar{z}-\rangle)$ |
| 2π | $-I$ | $- \bar{z}+\rangle$ |

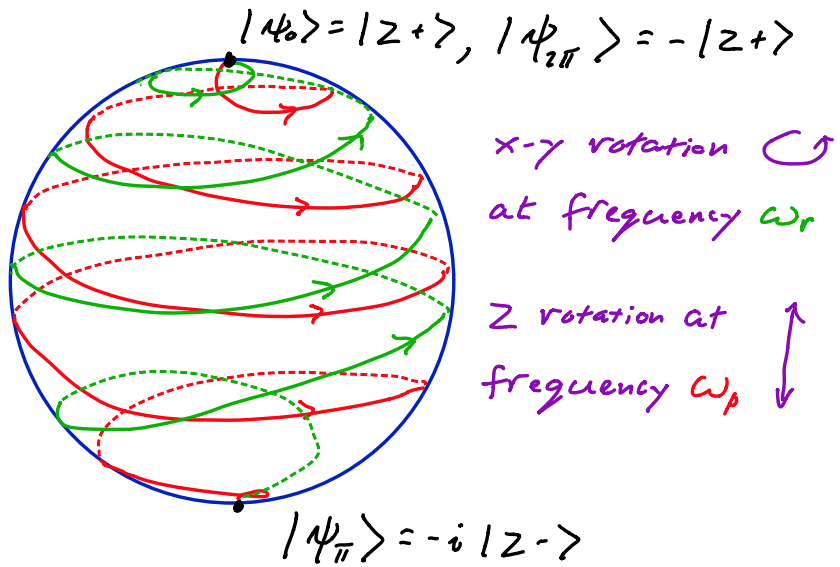
units $1/\omega_p$



Evolution of $\vec{S}(t)$

one full period $\frac{2\pi}{\omega_p}$

π pulse $t = \frac{\pi}{\omega_p}$



Off resonance Define $\delta = \omega_2 - \omega_r \Rightarrow \bar{H} = \frac{1}{2} \begin{pmatrix} \delta & \omega_p \\ \omega_p & -\delta \end{pmatrix}$

Eigenvalues $\lambda = \pm \frac{1}{2} \Omega, \Omega = \sqrt{\omega_p^2 + \delta^2}$

Eigenvectors $|\bar{\varphi}_{\pm}\rangle = A \begin{pmatrix} \delta \pm \Omega \\ -\omega_p \end{pmatrix} \quad A_{\pm} = \frac{1}{\sqrt{\omega_p^2 + (\delta \pm \Omega)^2}}$

$T = e^{-it\bar{H}} = e^{-i\Omega t/2} |\bar{\varphi}_+\rangle \langle \bar{\varphi}_+| + e^{i\Omega t/2} |\bar{\varphi}_-\rangle \langle \bar{\varphi}_-|$

Born Rule $P_r([\bar{z}-]_t | \psi_0 = |z+\rangle) = |\langle z- | T(t) | z+\rangle|^2$

$$= \frac{\omega_p^2}{\omega_p^2 + (\omega_r - \omega_2)^2} \text{Sin}^2\left(\frac{\Omega t}{2}\right)$$
