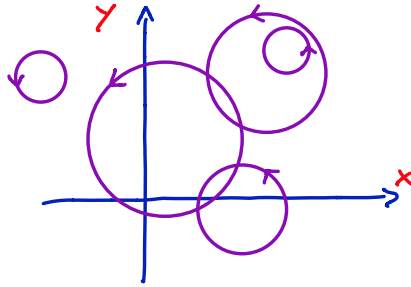


Charged particle in magnetic field  $\vec{B} = B\hat{e}_z = \vec{\nabla} \times \vec{A}$

Classical:  $\vec{f} = q\vec{v} \times \vec{B}$

Circle radius  $\sim \sqrt{E}$

"Cyclotron" frequency  $\omega_c = -qB/m$



Hamiltonian  $H = \frac{1}{2m} [\vec{p} - q\vec{A}]^2$  why?  $\begin{cases} \dot{\vec{r}} = \frac{\partial H}{\partial \vec{p}} \\ \dot{\vec{p}} = -\frac{\partial H}{\partial \vec{r}} = q\dot{\vec{r}} \times \vec{B} \end{cases}$

Gauge Invariance add  $\vec{\nabla} \chi(\vec{r})$  to  $\vec{A} \Rightarrow \vec{B}$  unchanged

But:  $H \rightarrow H' = \frac{1}{2m} [\vec{p} - q\vec{A} - q\vec{\nabla} \chi]^2$

Quantum case?  $H|\psi\rangle = E|\psi\rangle$  need  $H'|\psi'\rangle = E|\psi'\rangle$  same E  
 $\Rightarrow |\psi'\rangle = e^{iq\chi/\hbar} |\psi\rangle$  equivalent  $\psi$

Note: transformation indep. of  $E \Rightarrow$  holds for all quantum states

Coulomb gauge  $\vec{\nabla} \cdot \vec{A} = 0$   $\vec{A} = -\frac{1}{2} \vec{r} \times \vec{B} = \frac{1}{2} B (x\hat{e}_y - y\hat{e}_x)$

$$H = \frac{p_x^2 + p_y^2}{2m} - \underbrace{\frac{qB}{2m} (Xp_y - Yp_x)}_{\frac{\omega_c}{2} L_z} + \underbrace{\frac{q^2 B^2}{8m} (X^2 + Y^2)}_{\frac{1}{2} m \left(\frac{\omega_c}{2}\right)^2 (X^2 + Y^2)}$$

$$H = \underbrace{(N_d + N_g + 1)}_{2D \text{ H.O.}} \hbar \left(\frac{\omega_c}{2}\right) + \underbrace{(N_d - N_g)}_{L_z} \hbar \left(\frac{\omega_c}{2}\right) = (N_d + \frac{1}{2}) \hbar \omega_c$$

# States "Landau Levels"

$$|\Psi_{n_j, n_s}\rangle = \frac{1}{\sqrt{n_j! n_s!}} (a_d^\dagger)^{n_j} (a_g^\dagger)^{n_s} |\Psi_{00}\rangle$$

$\uparrow$  Control energy counter-clockwise  
 $\uparrow$  Control position clockwise

$$E_{n_j, n_s} = (n_j + \frac{1}{2}) \hbar \omega_c$$

$\uparrow$  degeneracy in  $n_g$

$$|\Psi_{00}\rangle = \frac{\kappa}{\sqrt{\pi}} e^{-\kappa^2(x^2 + y^2)/2} \quad \kappa^2 = qB/2\hbar$$

Alternate solution in Landau gauge  $\vec{A} = \vec{B} \times \hat{y}$   
 Breaks rotational symmetry

$$H = \frac{1}{2m} \{ P_x^2 + (P_y - qBx)^2 \}$$

translational symmetry in  $y$  not in  $x$

$\Rightarrow$  eigenstates  $\Psi(x, y) = e^{iky} \varphi(x)$

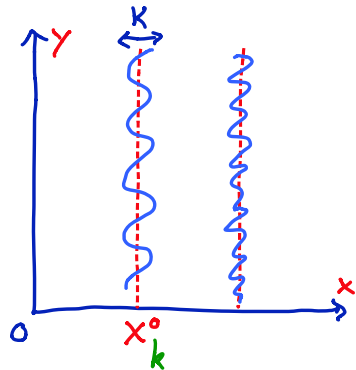
$$H\Psi = e^{iky} \frac{1}{2m} \{ P_x^2 + (\hbar k - qBx)^2 \} \varphi(x) = E e^{iky} \varphi(x)$$

$\therefore \varphi(x)$  is eigenstate of  $H_x = \frac{1}{2m} \{ P_x^2 + \frac{1}{2} m \omega_c^2 (x - \frac{\hbar k}{m\omega_c})^2 \}$

$\uparrow$  H.O. centered at  $x_k^0$   
 $\uparrow$   $x_k^0$

State  $\Psi_{nk}(x, y) = \frac{1}{\sqrt{n!}} (a^\dagger)^n |\varphi_0\rangle e^{iky}$

Energy  $E_n = (n + \frac{1}{2}) \hbar \omega_c$  independent of  $k$

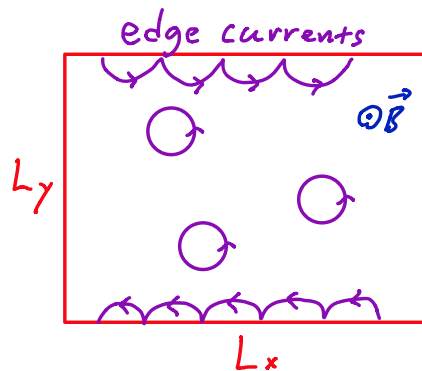


Degeneracy  $L_x \times L_y$  box with P.B.C.  $\Rightarrow k = \frac{2\pi}{L_y} N$

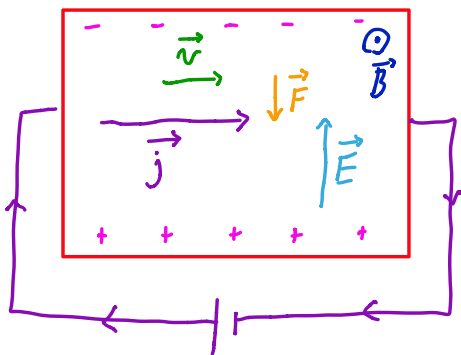
Box:  $0 \leq x < L_x \Rightarrow 0 \leq N < \frac{m\omega_c L_x L_y}{2\pi \hbar}$

$\therefore$  degeneracy  $D = \frac{qBA}{h} \equiv \Phi/\Phi_0$   $\Phi = BA$   $\Phi_0 = h/q$  "flux quantum"

1 Landau level / flux quantum



Integer quantum Hall effect



classical Hall resistivity  $\rho_{yx} = E_y/j_x = B/qn$  density

Density of states

