

Geometrical frustration in liquid Fe and Fe-based metallic glass

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We investigate short range order in liquid and supercooled liquid Fe and Fe-based metallic glass using *ab-initio* simulation methods. We analyze the data to quantify the degree of local icosahedral and polytetrahedral order and to understand the role of alloying in controlling the degree of geometric frustration. Comparing elemental Fe to Cu [1] we find that the degree of icosahedral order is greater in Fe than in Cu, possibly because icosahedral disclination line defects are more easily incorporated into BCC environments than FCC. In Fe-based metallic glass-forming alloys (FeB and FeZrB) we find that introducing small concentrations of small B atoms and large Zr atoms controls the frustration of local icosahedral order.

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I. INTRODUCTION

As noted by Frank [5], the local icosahedral clustering of 12 atoms about a sphere is energetically preferred because it is made up entirely of four-atom tetrahedra, the densest-packed cluster possible. However, local icosahedral order cannot be propagated throughout space without introducing defects. Frustration of packing icosahedra is relieved in

a curved space, where a perfect 12-coordinated icosahedral packing exists [6–8].

Disclination line defects of type $\pm 72^\circ$ may be introduced into this icosahedral crystal and thereby control the curvature. In order to “flatten” the structure and embed it in ordinary three dimensional space an excess of -72° disclinations is needed, and these cause increased coordination numbers of 14, 15 or 16. Large atoms, if present, would naturally assume high coordination numbers and aid in the formation of a disclination line network. Similarly, smaller atoms would naturally assume low coordination numbers of 8, 9 or 10, and have positive disclinations attached to them, increasing the frustration. For a particular coordination number, it may be possible to construct a cluster (known as a Kasper-Polyhedron [9]), made entirely of tetrahedrons, minimizing the number of disclinations.

Honeycutt and Andersen [10] introduced a method to count the number of tetrahedra surrounding an interatomic bond. This number is 5 for icosahedral order with no disclination, 6 for a -72° disclination and 4 for a $+72^\circ$ disclination. Steinhardt, Nelson and Ronchetti [11] introduced the orientational order parameter \hat{W}_6 to demonstrate short range icosahedral order. We employ both methods to analyze icosahedral order in supercooled metals and metal alloys, in addition to conventional radial distribution functions, structure factors and Voronoi analysis.

Many simulations have been performed on pure elemental metals and metal alloys using model potentials, [12–15], but do not necessarily produce reliable structures owing to their imperfect description of interatomic interactions. First principles (*ab-initio*) calculations achieve the most realistic possible structures, unhindered by the intrinsic inaccuracy of phenomenological potentials and with the ability to accurately capture the chemical natures of different elements and alloys. The trade-off for increased accuracy is a decrease in the system sizes one can study, so only local order can be observed, not long range. Also runs are limited to short time scales. Ab-initio studies on liquid Copper [16–18] and Iron [19] have not been analyzed from the perspective of icosahedral ordering. Recent ab-initio

studies on Ni and Zr [20, 21] find that the degree of icosahedral ordering increases with supercooling in Ni, while in Zr BCC is more favored. Studies on binary metal alloys by Jakse *et al.* [22, 23] and by Sheng *et al* [24] quantify local icosahedral order in the alloys. We previously [1] investigated icosahedral order in liquid and supercooled Cu

Elemental metals crystallize so easily that they can hardly be made amorphous at any given quench rate. Alloying can improve the ease of glass formation. For some special alloys, a bulk amorphous state can be reached by slow cooling. Pure elemental Fe is a poor glass former, but Fe-based compounds like FeB and especially FeZrB, show improved glass formability. We augment our molecular-dynamics simulation with another algorithm called ‘Tempering’ or ‘replica exchange method’ (REM) [36, 37] for fast equilibration at low temperatures.

In comparison to liquid and supercooled liquid Copper [1] which show only weak icosahedral order and very little temperature variation, Fe showed a monotonic increase in icosahedral order, which became very pronounced when supercooled. Analysis of quenched Fe revealed a natural way of introducing a single -72° disclination line segment into an otherwise perfect BCC environment, without disturbing the surrounding structure. Addition of B to Fe, decreased the icosahedral order, due to the positive disclinations centered on the smaller B atom, which increased frustration. Further inclusion of larger Zr atoms to form FeZrB found an enhanced icosahedral order compared to FeB. This could possibly be explained by formation of negative disclination line defects [28] anchored on the larger Zr atoms, which eases the frustration of icosahedral order on the Fe atoms.

At high temperatures all of our measured structural properties of liquid Cu [1] and liquid Fe resembled each other closely, and also strongly resembled a maximally random jammed [38] hard sphere configuration. This suggests that a nearly universal structure exists for systems whose energetics are dominated by repulsive central forces.

Section II describes our combined method of monte-carlo and first principles MD, that we refer to as “Tempering

MD” and discusses other simulation details. Section III presents our results on pure Fe while (section IV) compares this with FeB and FeZrB alloys.

II. TEMPERING MOLECULAR DYNAMICS (TMD) AND OTHER SIMULATION DETAILS

One reason alloys form glass more easily is that chemical identity introduces a new configurational degree of freedom that evolves slowly [39, 40]. Unfortunately, this makes simulation more difficult. It is especially difficult to equilibrate the system at very low temperatures, because the probability to cross an energy barrier drops, trapping it in particular configurations. For this reason we use a Monte-Carlo method, known as tempering or replica exchange [36, 37] to augment our first-principles MD, allowing us to sample the configurational space more efficiently than conventional MD.

In the canonical ensemble, energy fluctuates at fixed temperature. A given configuration C with energy E can occur at any temperature T with probability proportional to $e^{-\beta E}$, ($\beta = 1/kT$). Now consider a pair of configurations, C_1 and C_2 of energy E_1 and E_2 occurring in simulations at temperatures T_1 and T_2 . We can take C_1 as a member of the ensemble at T_2 , and C_2 as a member of the ensemble at T_1 , with a probability

$$P = e^{-(\beta_2 - \beta_1)(E_1 - E_2)} \quad (1)$$

without disturbing the temperature-dependent probability distributions of energy (or any other equilibrium property).

Because each run remains in equilibrium at all times even though its temperature changes, we effectively simulate a vanishingly low quench rate.

In practice we perform several MD simulations at temperatures separated by 100K. We use ultrasoft pseudopoten-

TABLE I: Details of tempering MD runs.

Chemical Species	Temperatures(K)	density (\AA^{-3})	time (ps)
Fe ₁₀₀	800-1900	0.0756	1.5
Fe ₈₀ B ₂₀	800-1500	0.0814	1.8
Fe ₇₀ Zr ₁₀ B ₂₀	800-1800	0.0787	1.8

tials [41] as provided with VASP [42] to perform the MD simulation. All calculations are Γ point calculations (a single k point). All runs use an MD time step of 2 fs, and reach total simulated time of order 1.5-1.8 ps (see Table I) with a total of $N=100$ atoms. Every 10 MD steps we compare the energies of configurations at adjacent temperatures and swap them with the above probability. Eventually, configurations initially frozen at low temperature reach a higher temperature. The simulations then can carry the structure over energy barriers, after which the temperature can again drop.

In an effort to explore the structures of compounds with differing glass-forming ability we compare pure elemental Iron and two Iron-based glass-forming alloys. Tempering MD requires that we perform simulations at a constant density for all the temperatures, but we have no rigorous means of predicting the density at high temperature. For pure liquid Iron, the density is known experimentally [43], and we use this value. For FeB and FeZrB, we took a high temperature liquid structure and quenched it, relaxing positions and cell lattice parameters, to predict a low temperature density. We then decreased the density of the relaxed structure by 6 percent to account for volume expansion, to arrive at the densities used in our liquid simulations.

Because of the efficient sampling of our tempering MD method, the structure of pure Fe partially crystallizes at low temperatures after about 1 ps. In the following discussion of our $T=800\text{K}$ sample we will refer to different structural

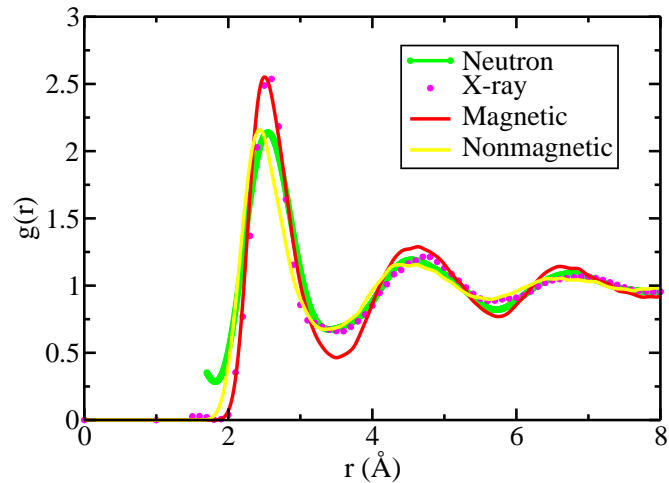


FIG. 1: Radial distribution function of pure elemental liquid Fe. Simulations (magnetic and non-magnetic) are run at $T=1800\text{K}$, compared with X-ray experiments at $T=1833\text{K}$ [43] and Neutron scattering experiments at 1830K [27].

features before and after crystallization. We also performed several long (2.0ps) conventional first-principles MD at $T=800\text{K}$ yielding results similar to the results of tempering MD prior to crystallization.

For all runs we employed spin polarization, reasoning that local magnetic moments exist even above the Curie point. These local moments have a significant influence on the short-range order because ferromagnetic Iron prefers a longer bond length than paramagnetic Iron [44]. Of course, the ferromagnetic state of the liquid implies improper long-range correlations. Unfortunately, since our forces are calculated for electronic ground states, we cannot rigorously model the true paramagnetic state of liquid Iron and Iron-based alloys with these methods.

III. PURE FE

A. Radial Distribution Function $g(r)$

The radial distribution function, $g(r)$, is proportional to the density of atoms at a distance r from another atom and is calculated here by forming a histogram of bond lengths. We use the repeated image method to obtain the bond lengths greater than half the box size and anticipate $g(r)$ in this range may be influenced by finite size effects. Further, we smooth out the histogram with a gaussian of standard deviation 0.05\AA .

To evaluate the role of magnetism, Fig. 1 illustrates the radial distribution functions for liquid Fe simulated at $T=1800\text{K}$, just below melting ($T_m=1833\text{K}$). Evidently, the simulation with magnetism yields good agreement in the position and height of the first peak in $g(r)$ with experimental X-ray $g(r)$ [43], while neglect of magnetic moments results in near neighbor bonds that are too short and too weak. However, magnetism overestimates the strength of long-range correlations beyond the nearest-neighbor peak, while neglecting magnetism yields reasonably accurate long-range $g(r)$. Nevertheless, for the present study of local order, it is necessary to make spin polarized calculations, to get the short range correlations and hence the local order correct. Strangely, a recent experimental neutron $g(r)$ [27] has a shorter and broader first peak compared to both our magnetic $g(r)$ and to the $g(r)$ from the prior X-ray diffraction experiment [43]. The positions of the different maxima and minima in our simulated magnetic $g(r)$ compare well with both the experiments. The position of the first peak in our magnetic $g(r)$ is shifted by 0.05\AA to the left of the neutron first peak. The X-ray experiment doesn't have enough data points around the first maximum to determine the peak position accurately.

We calculate the coordination number by counting the number of atoms within a cutoff distance from a central

atom. We choose the cutoff distance (R_{cut}) at the first minimum of $g(r)$. For pure Fe the minimum is at $R_{cut}=3.5 \text{ \AA}$. The precise location of the minimum is difficult to locate, and its variation with temperature is smaller than the error in locating its position, so that we don't change the value of R_{cut} with temperature. With this value of R_{cut} we find an average coordination number (N_c) of 13.2 which is nearly independent of temperature (N_c changes from 13.1 at high temperature to 13.3 with supercooling).

B. Liquid Structure Factor $S(q)$

The liquid structure factor $S(q)$ is related to the radial distribution function $g(r)$ of a liquid with density ρ by,

$$S(q) = 1 + 4\pi\rho \int_0^{\infty} [g(r) - 1] \frac{\sin(qr)}{qr} r^2 dr \quad (2)$$

One needs the radial distribution function up to large values of r to get a good $S(q)$. In our first principles simulation, we are restricted to small values of r , due to our small system sizes, so we need a method to get $S(q)$ from our limited $g(r)$ function. Baxter developed a method [47, 48] to extend $g(r)$ beyond the size of the simulation cell. The method exploits the short range nature of the direct correlation function $c(r)$, which has a range similar to the interatomic interactions [49], as opposed to $g(r)$ which is long ranged.

Assuming that $c(r)$ vanishes beyond a certain cutoff distance r_c , we solve the Baxter's equations iteratively to obtain the full direct correlation function for $0 < r < r_c$. From $c(r)$ we calculate the structure factor $S(q)$ by a standard Fourier Transform. The $S(q)$ showed good convergence with different choices of r_c . A choice of $r_c=5\text{\AA}$ seemed appropriate because it was one half of our smallest simulation cell edge length. Even though in metals there are long range oscillatory Friedel oscillations, our ability to truncate $c(r)$ at $r_c=5\text{\AA}$ shows that these are weak compared with

short range interactions. An application of this method to obtain $S(q)$ of Cu [1] showed excellent agreement with the experimental $S(q)$.

The simulated $S(q)$ for pure Fe at $T=1800\text{K}$ (see Fig. 2) is compared to recent neutron scattering experiments at $T=1830\text{K}$. Even though the positions of the different peaks compare very well, there is serious discrepancy in their heights. Especially, the height of the first peak of our simulated $S(q)$ is higher than that of the experiments. This discrepancy is expected because we include magnetism, which gives accurate short-range correlations while overestimating the long-range ones (see Fig. 1). But the cause of the discrepancy is not entirely clear since a comparison of the simulated structure factor of Ni [20] (done without including magnetism) with neutron scattering experiments [27] shows similar discrepancies between their $S(q)$'s.

A sum rule can be obtained for $S(q)$ [45, 46]. By inverting the fourier transform of Eq. 2 and then taking the $r \rightarrow 0$ limit, one gets

$$I(Q) \equiv \int_0^Q q^2 [S(q) - 1] dq \rightarrow -2\pi^2 \rho \quad (3)$$

in the limit $Q \rightarrow \infty$. Further, the integral is supposed to oscillate with Q about the limiting value as $Q \rightarrow \infty$. Using our $S(q)$ we observed that the integral is consistent with the sum rule and oscillates nicely about the limiting value for $Q \geq 3\text{\AA}^{-1}$, while using the $S(q)$ from the neutron scattering experiments [27], we observe a positive drift in the mean value about which the integral oscillates. Such a drift could indicate the presence of spurious background corrections. The $S(q)$ from the X-ray experiment [43] seems to be in good agreement with the ideal sum rule.

As we lower the temperature, the peak heights in $S(q)$ grow, indicating an increase in short range order with supercooling. We also observe a slight shoulder in the second peak of the $S(q)$ (Fig. 2), which grows with supercooling.

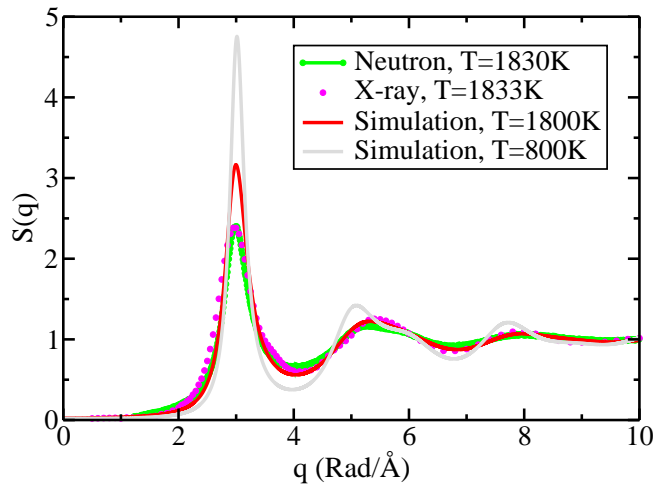


FIG. 2: Comparison of simulated and experimental $S(q)$ near the melting temperature of Fe.

TABLE II: \hat{W}_6 values for a few clusters

Cluster	HCP	FCC	ICOS	BCC	Z14	Z15	Z16	Z8	Z9	Z10
No. of atoms	12	12	12	14	14	15	16	8	9	10
\hat{W}_6	-0.012	-0.013	-0.169	+0.013	-0.093	-0.037	doesn't exist	+0.01	-0.038	-0.093

The split second peak positions are in the ratios of 20:12 and 24:12 with respect to the first peak positions, just what one would ideally observe if there was icosahedral order [28, 29].

C. Bond Orientation Order Parameter \hat{W}_6

Steinhardt, *et al.* [11] introduced the \hat{W}_l parameters as a measure of the local orientational order in liquids and undercooled liquids. To calculate \hat{W}_l , the orientations of bonds from an atom to its neighboring atoms are projected onto a basis of spherical harmonics. Rotationally invariant combinations of coefficients are then averaged over many atoms in an ensemble of configurations. The resulting measures of local orientational order can be used as order

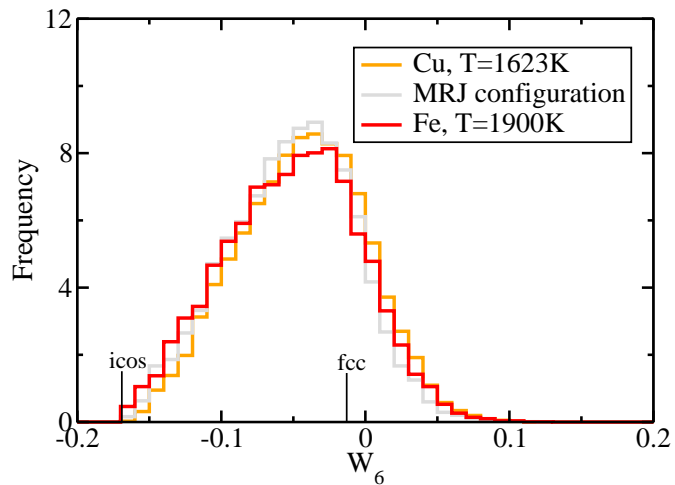


FIG. 3: Distribution of \hat{W}_6 in liquid and supercooled liquid Fe and Cu. Fe shows more pronounced icosahedral order than Cu with supercooling

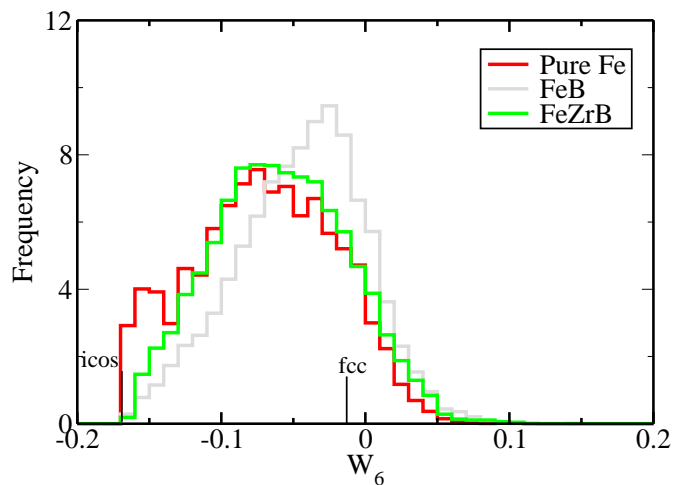


FIG. 4: Distribution of \hat{W}_6 around Fe atoms, in supercooled Fe, FeB and FeZrB at $T=800\text{K}$

parameters to characterize the liquid structures. For an ideal icosahedral cluster, $l = 6$ is the minimum value of l for which $\hat{W}_6 \neq 0$. Table II enumerates \hat{W}_6 values for different ideal clusters. The ideal icosahedral value of \hat{W}_6 is far from other clusters, making it a good icosahedral order indicator. The clusters Z14 to Z10 are deltahedrons, with their surface bonds triangulated with equilateral triangles (except for Z16 for which this is not possible). As one

starts including 72° disclinations to the icosahedron, which has zero disclinations, keeping the resulting cluster to be a deltahedron, one finds that the shift in \hat{W}_6 is very nearly the same for both the positive and the negative types of the disclination. There doesn't seem to be a strict deltahedron present for the Z16 cluster. But it is possible to create it to be nearly a perfect deltahedron in which case its \hat{W}_6 value is very close to that of the Z8.

As before, we choose the cutoff distance to specify near neighbors as $R_{cut} = 3.5\text{\AA}$. For pure Fe at high temperatures, \hat{W}_6 resembles that of high temperature liquid Cu [1], which in turn resembles the maximally random jammed configuration of hard spheres [38](Fig. 3). On this basis we suggest that the MRJ configuration represents an idealized structure that is universal for strongly repulsive interactions. All pure metallic systems might approach this ideal structure at sufficiently high temperature.

However, as temperature drops the \hat{W}_6 distribution shifts strongly to the left, with a pronounced increase in $\hat{W}_6 \leq -0.1$ (Fig. 4). This indicates a rather high concentration of nearly icosahedral clusters in supercooled liquid Fe. Even at very low percentages of supercooling of Cu ($\sim 3\%$ [1]) and Fe ($\sim 2\%$, not shown), Fe shows a clear enhancement in the negative \hat{W}_6 distribution as compared to Cu.

We checked the \hat{W}_6 distribution of a non-magnetic simulation to see if it is strongly influenced by magnetism, and found a nearly identical result. In particular, we still found a strong enhancement of the nearly icosahedral clusters relative to liquid Cu or pure Fe at high temperature.

D. Voronoi Analysis

To explain the origin of this low \hat{W}_6 peak in pure Fe, we performed a Voronoi analysis [50] of the liquid before and during crystallization. A Voronoi polyhedron is described by indices (F_3, F_4, F_5, \dots) where F_i denotes the number of

TABLE III: Voronoi statistics (% of Fe atoms) for typical Fe-centered clusters in Fe, FeB and FeZrB. Edges smaller than 0.33\AA and faces smaller than 0.3\AA^2 have been treated as a single vertex

Voronoi type	Fe				FeB	FeZrB
	supercooled		defective		supercooled	supercooled
	liquid		crystal		liquid	liquid
	800K	relaxed	800K	relaxed	800K	800K
(0,0,12)	8.3	11.0	5.0	0.0	0.63	1.4
(0,0,12,2)	2.1	3.6	2.0	2.0	0.5	0.3
(0,2,10,{0,1})	6.2	6.8	2.0	0.0	4.4	7.1
(0,1,10,2)	3.6	5.4	3.0	0.0	2.1	3.1
(0,4,8,{0,1,2,3})	7.9	4.8	10.0	6.0	8.6	9.4
(0,3,8,{0,1,2,3})	6.8	6.0	12.0	0.0	4.3	8.3
(0,2,8,{1,2,3,4})	3.0	3.0	3.0	0.0	0.0	0.0
(0,3,6,4)	2.1	2.2	2.0	6.0	0.0	0.0
(0,5,4,4)	1.5	1.2	3.0	18.0	0.0	0.0
(0,6,0,8)	0.0	0.4	0.0	56.0	0.0	0.0

faces with i edges. For example (0,0,12) denotes an icosahedron, while (0,0,12,2) denotes a 14-coordinated (Z14) atom, with 12 5-fold bonds and 2 6-fold bonds. The (0,0,12,2) is a characteristic TCP (tetragonal close-packed) structure of the Frank-Kasper type, with a -72° disclination line running through an otherwise perfect icosahedron. In a body-centered cubic crystal all atoms are of Voronoi type (0,6,0,8), which is an alternate 14-coordinated structure.

Supercooled liquid Fe at $T=800\text{K}$ (Table III) contains a high fraction of icosahedral atoms of type (0,0,12) and

(0,0,12,2) . Those with Voronoi type (0,2,10,{0,1}) and (0,1,10,2) also have very negative \hat{W}_6 , so that they can be thought of as related to the icosahedron. Together, they explain the enhanced negative \hat{W}_6 distribution. Supercooled Cu (which shows weak icosahedral order) and high temperature Fe (T=1900K), in contrast contain about 1.2% of (0,0,12) and no (0,0,12,2). The MRJ configuration contains 1.2% (0,0,12) but no (0,0,12,2).

Strikingly, the icosahedral clusters tend to join in pairs and strings of 3 atoms in length. An instantaneous quench of a particular liquid structure at 800K which had a high fraction of icosahedral units, using conjugate gradient relaxation of the atomic coordinates and lattice parameters, shows a clear enhancement of the icosahedral and other closely related units. The strings of icosahedral units found in the supercooled liquid connected to form networks. A quench starting from a different instantaneous supercooled liquid structure at 800K containing fewer icosahedral structures resulted in rapid crystallization. This indicates that the presence of icosahedrons may inhibit crystallization. Similar quenches, starting from instantaneous liquid structures at higher temperatures, also resulted in relaxed structures with some of them showing a high negative \hat{W}_6 distribution comparable to the distribution at 800K. The quenches that partially crystallized showed a high fraction of BCC (0,6,0,8)'s ($\sim 40\%$) and was always accompanied by a high fraction of (0,5,4,4)'s ($\sim 25\%$), which are otherwise absent or very low in the liquid.

Under TMD the supercooled liquid at T=800K eventually crystallizes, with the \hat{W}_6 distribution peaked strongly around zero consistent with the value for ordinary crystalline clusters, but retaining a subset of atoms with nearly icosahedral $\hat{W}_6 < -0.14$. Not surprisingly, these were precisely the atoms that had Voronoi type (0,0,12) prior to crystallization. The Z14 (0,0,12,2) atoms (with $\hat{W}_6 = -0.07$) were mutual near-neighbors, linked along their 6-fold (-72°) bonds, and also were neighbors of the nearly icosahedral atoms.

We quenched this sample by conjugate gradient relaxation of atomic coordinates and lattice parameters. After

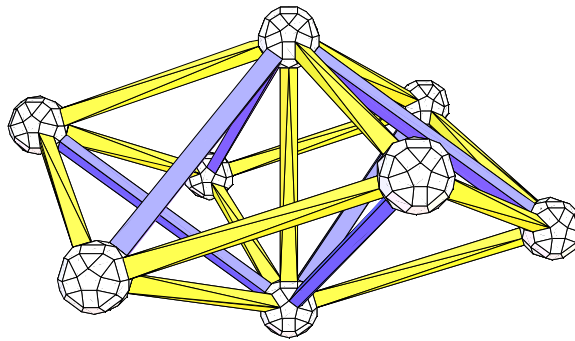


FIG. 5: Fragment of BCC crystal illustrating a central bond in the $[1,1,1]$ (vertical) direction, surrounded by two equilateral triangles rotated 60° from each other and displaced on either side of the midplane. Balls and sticks follow ZomeTool convention: Balls are icosahedral, yellow sticks are 3-fold and blue sticks are 2-fold bonds.

relaxation we found 56 atoms had BCC Voronoi type $(0,6,0,8)$. Six icosahedral atoms became 12-coordinated $(0,4,8)$ structures surrounding the bond connecting the two $(0,0,12,2)$ disclinated icosahedral atoms, which retained their type. The remaining atoms served to link the cluster of icosahedron-related atoms to the surrounding defect-free BCC crystal.

The icosahedron-related structure thus forms a point defect in an otherwise perfect BCC crystal. A simple way to create this defect is to take two consecutive triangles surrounding a near-neighbor bond along the BCC $[1,1,1]$ direction (see Fig. 5), displace them into the perpendicular bisecting plane and rotate by 30° . The atoms along the $[1,1,1]$ bond are now connected to each other by a six-fold bond (-72° disclination) and are connected to the six displaced atoms by five-fold bonds. This scheme to transform a BCC rhombic dodecahedron into a Frank-Kasper Z14 polyhedron was described in Ref [51, 52] in an attempt to explain certain diffraction anomalies in iron and vanadium-based alloys under ion irradiation.

Manually removing the defect, by reversing the above procedure then relaxing, yields a perfect BCC crystal with

all atoms in a (0,6,0,8) Voronoi environment. We also embedded the point defect in an otherwise perfect BCC crystal with 128 atoms in a cubic box. Relaxation showed that the defect was stable. The energy of the defect was 6.18eV and the fractional volume increase was 0.013.

Since tungsten, like Iron, crystallizes in BCC, we performed a separate first-principles simulation of liquid Tungsten. After supercooling by 13% to $T=3200\text{K}$, we found the \hat{W}_6 histogram resembled that of Fe at 1600K (which is 15% supercooled). A Voronoi analysis revealed a similar percentage of (0,0,12)'s and (0,0,12,2)'s in both these elemental BCC-forming metals. Hence we believe the relationship between BCC and icosahedral structures may be linked to our observed high fraction of icosahedra in liquid Fe as compared to Cu. This may also explain the reason why we are able to supercool BCC Fe more deeply than FCC Cu in our simulations.

We also included an icosahedron point defect inside an otherwise perfect FCC crystal of Cu, with 256 atoms in a cubic box. Relaxation showed that the defect was stable. The energy cost of the defect was 4.69eV and the fractional volume increase was 0.012. Even though this FCC defect cost less energy than the point defect in BCC, and also needs little rearrangement of atoms as reflected by the slightly lower fractional density change, we do not see a significant icosahedral order in liquid Cu, when compared to liquid Fe. This could be because the number of icosahedrons that a single Z14 disclination can stabilize (up to six) is greater than one. So even an equal number of the two different defects in BCC and FCC would result in more icosahedral order in BCC than in FCC.

E. Honeycutt and Andersen analysis

Honeycutt and Andersen [10] introduced a useful assessment of local structure surrounding interatomic bonds. We employ a simplified form of their analysis, counting the number of common neighbors shared by a pair of near-neighbor

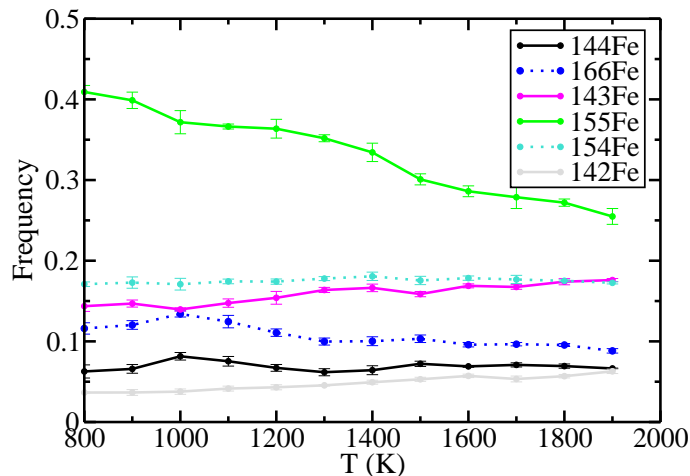


FIG. 6: Honeycutt-Andersen analysis for pure Fe shows a clear increase in five-fold bonds with supercooling.

atoms. This identifies the number of atoms surrounding the near-neighbor bond and usually equals the number of edge-sharing tetrahedra whose common edge is the near-neighbor bond. We assign a set of three indices to each bond. The first index is 1 if the root pair is bonded (separation less than or equal to R_{cut}). The second index is the number of near-neighbor atoms common to the root pair, and the third index gives the number of near-neighbor bonds between these common neighbors. We take the same value of $R_{cut}=3.5\text{\AA}$ as mentioned before. Note that the Honeycutt and Andersen fractions depend sensitively on R_{cut} , making precise quantitative comparisons with other prior studies difficult.

In general, 142's are characteristic of close packed structures (FCC and HCP) and 143's are characteristic of distorted icosahedra [53]. They can also be considered as $+72^\circ$ disclinations [6–8]. Likewise, 15's are characteristic of icosahedra, with 155's characterizing perfect icosahedra while 154's and 143's characterize distorted ones. 16's indicate -72° disclinations. 166's and 144's are also characteristic of BCC.

A Honeycutt and Andersen analysis for pure Fe, with an $R_{cut}=3.5\text{\AA}$ (Fig. 6), showed that with supercooling the

TABLE IV: Average coordination number in FeB.

	FeFe	FeB	BFe	BB
R_{cut}	3.4	3.0	3.0	2.3
$N_{\alpha\beta}$	12.0	2.2	8.9	0.4

fraction of 15 bonds rises from 0.46 at T=1900K to 0.59 at T=800K in the liquid before crystallization. The fraction of 155's (characteristic of perfect icosahedra) was always larger than 154's (characteristic of distorted icosahedra), and seemed to be steeply increasing with supercooling as opposed to 154's, which were relatively flat. The fraction of 14 bonds drop from 0.32 to 0.30, with the icosahedral 143's being always higher than the cubic 142's. The 144's, which are characteristic of BCC remain nearly flat, even though the 166's show a slight increase. The ease of embedding a Z14 disclination in BCC Fe (see section III D) might explain the slight increase in the 166's.

IV. FE-B AND FE-ZR-B

A. Radial Distribution Function $g(r)$

Fig. 7 shows the pair correlation functions of supercooled FeZrB at T=800K. From the heights of the first peaks, we see that the strongest bonds form between the metalloid (B) and the metal (Fe or Zr). The relative bond lengths reveal, as expected, that B behaves as a small atom, Fe is medium sized and Zr is large. Comparing simulations of FeZrB to pure Fe, we find that g_{FeFe} is reduced by alloying with B or ZrB, because Fe prefers to associate with B or Zr rather than with Fe.

To calculate the coordination number, define $N_{\alpha\beta}$ as the average number of atoms of type β around an atom of type

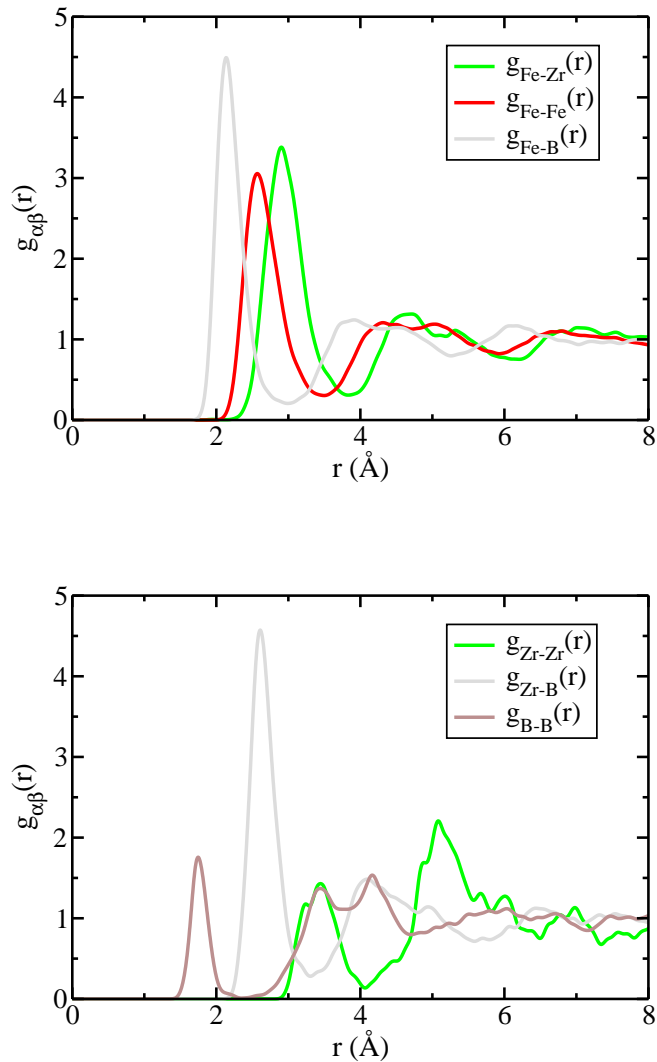


FIG. 7: Partial radial distribution functions in FeZrB at T=800K.

TABLE V: Average coordination number in FeZrB.

	FeFe	FeZr	FeB	ZrFe	ZrZr	ZrB	BFe	BZr	BB
R_{cut}	3.4	3.8	3.0	3.8	4.0	3.3	3.0	3.3	2.4
$N_{\alpha\beta}$	9.5	2.0	2.0	14.0	0.9	2.8	6.8	1.4	0.4

α , We set R_{cut} at the first minima of the partial radial distribution functions (Fig. 7). We list the partial coordination numbers of FeB in Table. IV and FeZrB in Table. V (averaged over all the temperatures since the temperature

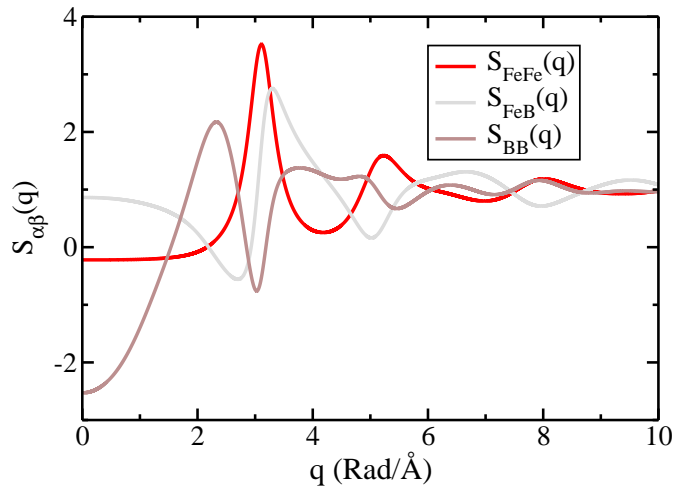


FIG. 8: Simulated partial structure factors of FeB at T=800K

dependence is very weak and non-monotonic). The average value of N_{FeFe} decreases with alloying, due to decrease in the concentration of Fe and the preference to bind with B and Zr. Zr, being a large atom, has a larger coordination number. Also note that B and Zr favor each other more than themselves. We do find some B-B pairs in the liquid state.

B. Liquid Structure Factor $S(q)$

Fig. 8 shows the Faber-Ziman [54] partial structure factors of FeB at the lowest simulated temperature of T=800K, defined as,

$$S_{\alpha\beta}(q) = 1 + 4\pi\rho \int_0^{\infty} [g_{\alpha\beta}(r) - 1] \frac{\sin(qr)}{qr} dr \quad (4)$$

The positions of the first and second peaks in partial $S_{FeFe}(q)$ is in very good agreement with the experimental results for amorphous FeB [55]. Also, at the position of the splitting of the second peak in $S_{FeFe}(q)$, as observed in the experiments, we observe a slight shoulder. Similarly, the positions of the different peaks in $S_{FeB}(q)$ and $S_{BB}(q)$ are

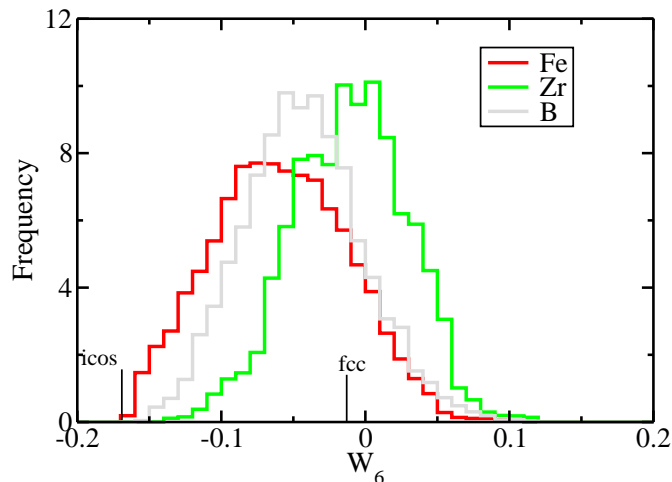


FIG. 9: Distribution of \hat{W}_6 for different chemical species in supercooled FeZrB at T=800K

in agreement with the experiments.

For FeB, the $q \rightarrow 0$ limit of $S_{FeFe}(q)$ is comparable to the experimental value. The $q \rightarrow 0$ limit in the other two partials differ from the experiment, $S_{BB}(q)$ more seriously than $S_{FeB}(q)$. We think that this discrepancy in the long wavelength regime is probably due to the very low density of B in our system leading to poor statistics. Nevertheless, the excellent agreement in the positions of the different peaks in the partial structure factors shows that we have reasonably good representative structures of FeB at T=800K. Partial structure factors or partial pair distribution functions are not available experimentally to compare with FeZrB simulations.

C. Bond Orientation Order Parameter \hat{W}_6

To define the \hat{W}_6 distribution in an alloy, we concentrate on central atoms of some particular species (e.g. Fe) but consider the neighboring atoms of all species. We chose the near-neighbor cutoff distances as before. Compare the Fe-based \hat{W}_6 distributions at their supercooled temperatures in Fig. 4. The origin of high negative \hat{W}_6 values for pure

Fe was previously explained in section III D. Replacing a few medium sized Fe atoms with smaller B atoms causes negative disclination lines to concentrate on Fe, leading to a drop in the ideal icosahedral clustering on these Fe atoms and strongly reducing the extreme negative values of \hat{W}_6 in FeB. On the contrary, inclusion of Zr in FeZrB causes negative disclinations to attach to Zr, easing frustration, leading to more Fe centered clusters with icosahedral ordering and increasing the negative region of \hat{W}_6 , as compared to FeB. Inclusion of the big Zr atom enhances icosahedral order on the Fe atoms.

Fig 9 shows the \hat{W}_6 distributions for FeZrB with centers at Fe, Zr and B. The histogram with the center at Zr is almost symmetric about the value of zero. This suggests that the local environment about Zr atoms is nearly spherical, as is expected given its large size. The B centered \hat{W}_6 histogram is also asymmetric towards negative \hat{W}_6 values due to Kasper-Polyhedrons that are deltahedrons with maximum five fold bonds for a particular coordination and slightly distorted versions of them.

D. Voronoi Analysis

A Voronoi analysis was performed for FeB in the supercooled liquid at T=800K (Table III). The Fe environments were mostly (0,4,8,x) or (0,3,8,x) types, where $x=\{0,1,2,3,4\}$, with the higher coordination polyhedron being more favored. This was followed by the very negative \hat{W}_6 valued (0,0,12), (0,2,10,x)'s and (0,1,10,2)'s Voronoi types occurring at lower frequency than in pure Fe. Boron mainly had environments of type (0,3,6)'s ($\sim 15\%$ of B atoms) (Kasper polyhedron for Z=9 containing a $+72^\circ$ disclination) and (0,5,4)'s ($\sim 20\%$) typical of the tri-capped trigonal prism (TTP) and the mono-capped square archimedean prism (slightly distorted variant of TTP) respectively, with the (0,5,4)' being more predominant than the TTP structure. The TTP is found in the crystal structures of Fe₃B with

Pearson symbols oP16 and tI32. The distorted (0,5,4) version is found in the structure Fe_{23}B_6 of Pearson type cF116. These structures have been identified as the leading competitors for B-Fe glass [56]. Boron also took environments of the Kasper polyhedron (0,4,4) ($\sim 3\%$) corresponding to $Z=8$ and (0,2,8) ($\sim 10\%$) corresponding to $Z=10$. The association of B with $+72^\circ$ disclinations explains how it increases the frustration of icosahedral order on the Fe atoms. Clearly the improved glass-formability of FeB compared with elemental Fe cannot be due to icosahedral order. Rather, it is presumed to be caused by the deep eutectic at $\text{Fe}_{83}\text{B}_{17}$.

A Voronoi analysis of supercooled FeZrB at $T=800\text{K}$ shows a clear increase in the very negative \hat{W}_6 polyhedrons, and also a decrease in the number of Z14 (0,0,12,2) types on Fe atoms indicating a decrease in frustration in the ternary as compared to the binary. Environments around B atoms were roughly similar in the binary and the ternary, with a slight increase in the lower coordinated ($Z=8$) Kasper polyhedrons at the cost of higher coordinated ($Z=10$) ones. Zirconium took a variety of polyhedra, with an average coordination of 17.6 and a minimal coordination of 15, owing to its large size compared to the other constituents.

E. Honeycutt and Andersen analysis

We made a HA analysis of the ternary glassy alloy, by looking at root pairs of chemical species α and β , choosing R_{cut} in the manner of the \hat{W}_6 analysis. The frequency is normalized to sum to one for each species pair $\alpha\beta$. Table VI lists the fraction of different $1x$ pairs in supercooled FeZrB.

Among the $1x$ pairs with Fe as one of the root pairs, the 15's are most abundant at all temperatures. The 15's are mainly comprised of 155's and the 154's, with the 155's being always higher than the 154's. The percentage of 15's is similar for FeZrB, FeB as well as pure Fe. Note that the 15's are largest for the Fe-Fe pairs. They also show

TABLE VI: HA analysis for supercooled FeZrB at T=800K

Root Pair						
	Fe-Fe	Fe-Zr	Fe-B	Zr-Zr	Zr-B	B-B
14 pairs	0.30	0.19	0.54	0.11	0.55	0.75
142 pairs	0.08	0.06	0.03	0.06	0.03	0.02
143 pairs	0.17	0.11	0.35	0.04	0.33	0.33
144 pairs	0.05	0.02	0.17	0.01	0.19	0.41
15 pairs	0.56	0.53	0.39	0.48	0.41	0.18
154 pairs	0.22	0.27	0.03	0.36	0.04	0.0
155 pairs	0.33	0.22	0.37	0.08	0.36	0.17
16 pairs	0.10	0.24	0.01	0.36	0.01	0.01
166 pairs	0.09	0.20	0.0	0.26	0.00	0.00

a steady enhancement with supercooling, unlike the 14's which decrease with supercooling. Among the 14's, the icosahedron-related 143's for all root pairs are always higher than the close-packed 142's, and remain fairly constant with supercooling. The 14's are maximal for FeB root pairs and minimal for FeZr.

Supercooled pure Fe has a high percentage of "16"'s (13%) compared to pure Cu (7%). The high number of "16"'s in Fe is related to the occurrence of Z14 (0,0,12,2) environments in which the 6-fold bond carries a -72° disclination, rather than the BCC 14 atom arrangement which would also show a high degree of four-fold and six-fold bonds (section IIID). Adding B to Fe, shows an increase of "16"'s for Fe-Fe pairs to 21% in FeB (not shown). Adding

large Zr atoms to FeB decreases the occurrence of 16's on the Fe-Fe pairs, putting them on Fe-Zr pairs and easing the frustration of Fe centers. This causes the geometry about Fe centers to be more icosahedral, and hence the shift of the \hat{W}_6 towards negative values.

Among the root pairs not containing Fe, note that the B-B pair has the maximum 14's ($+72^\circ$ disclinations), especially 144's, while the Zr-Zr pair has the maximum 16's (-72° disclinations), especially the 166's, emphasizing the role of size in controlling the frustration in alloys.

V. CONCLUSION

This study quantifies icosahedral and polytetrahedral order in supercooled liquid metals and alloys. This is the first such analysis of glass-forming Fe compounds using configurations from first-principles simulations. While the structural properties of Fe and Cu strongly resemble each other at high temperature, and indeed are close to a maximally random jammed structure [1, 38], their behavior evolves substantially, and in different manners, as the liquid is supercooled. Proper modeling of atomic interactions is essential to capture the differing behavior of each element, and use of a first-principles simulation is the most reliable means of achieving this.

For pure elements we find the degree of local icosahedral order in the supercooled liquid depends on the low temperature crystal structure, with BCC metals such as Fe and W accommodating icosahedra more readily than the FCC element Cu. Alloying with large or small atoms can further influence the degree of icosahedral order, with small atoms (e.g. B in Fe) aggravating the frustration by introducing positive disclination line defects, while large atoms (e.g. Zr in Fe) naturally stabilize negative disclination line defects, relieving frustration on the medium-sized Fe atoms. While the enhanced glass-forming ability of FeB compared to Fe cannot be related to icosahedral order,

we suggest that the formation of icosahedral order and disclination line network, together with the slow dynamics of chemical ordering in a complex alloy, enhances the glass-forming ability of FeZrB compared with FeB.

Acknowledgments

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